

On the time-slice axiom in 2d conformal AQFT

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Online talk @ York, 20 January 2022.

Joint work w/ M. Benini and L. Giorgetti [arXiv:2111.01837]

Orthogonal categories and AQFTs

- Def: An orthogonal category $\overline{G} = (G, \perp)$ is a category G with a symmetric and o-stable subset

$$\perp = \left\{ M_1 \xrightarrow{f_1} M' \xleftarrow{f_2} M_2 \right\} \subseteq \text{Mor } G \times_{\text{t}} \text{Mor } G.$$

An AQFT on \overline{G} is a \perp -commutative functor $A: G \rightarrow \text{Alg} \subseteq {}^* \text{Alg}_{\text{As}}(\mathbf{T})$.

\uparrow
ISMCat

$$\begin{array}{ccc}
 A(M_1) \otimes A(M_2) & \xrightarrow{A(f_1) \otimes A(f_2)} & A(M')^{\otimes 2} \\
 \downarrow & & \downarrow \mu_{M'} \\
 A(M_1) \otimes A(M_2) & \xrightarrow{\quad \text{if } f_1 \perp f_2 \quad} & A(M') \\
 & & \downarrow \\
 A(M')^{\otimes 2} & \xrightarrow{\quad \text{op} \quad} & A(M')
 \end{array}$$

- Ex: Loc_m w/ \perp = causally disjoint

and m-dim. LCQFTs satisfying Einstein causality
but NOT necessarily time-slice

Nice features of this formalism

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Prop: The assignment $\bar{G} \mapsto AQFT(\bar{G})$ of AQFT categories can be upgraded to a 2-functor $AQFT: OrthCat^{op} \longrightarrow \mathbf{Cat}$.

Prop: Denote by $AQFT(\bar{G})^W \subseteq AQFT(\bar{G})$ the full subcategory of AQFTs satisfying time-slice for a set of maps $W \subseteq \text{Mor } G$, i.e.

$A(f): A(M) \xrightarrow{\sim} A(M')$ is iso for all $(f: M \rightarrow M') \in W$.

The localization functor $L: \bar{G} \rightarrow \bar{G}[W^{-1}] = (G[W^{-1}], L_{\bar{G}}(I_{\bar{G}}))$ induces an equivalence of categories

$$L^*: AQFT(\bar{G}[W^{-1}]) \xrightarrow{\sim} AQFT(\bar{G})^W.$$

In simple words: The time-slice axiom can be "solved" by computing the localized category $G[W^{-1}]$.

Known cases where localizations have been computed

1.) Haag-Kastler type AQFTs on M [Benini, Dappiaggi, AS]

$$C\text{Open}(M)[\text{Cauchy}^{-1}] \simeq (\text{Open}(M))^D = \left\{ U \subseteq M : \begin{array}{l} DU = U \\ \uparrow \\ \text{Cauchy development} \end{array} \right\}$$

2.) The "RCE category" [Bruinsma, Fewster, AS]

$$\left(\begin{array}{ccc} & M_+ & \\ M & \xleftarrow{R} & \xrightarrow{M_h} \\ & M_- & \end{array} \right) [\text{Cauchy}^{-1}] \simeq B\mathbb{Z} = \begin{array}{c} \mathbb{Z} \\ \circlearrowleft \end{array}$$

the generator $1 \in \mathbb{Z}$ gives $\text{rce}_{(M,h)}$!

3.) 1d LCQFT [Bruinsma or AS, don't remember]

$$\text{AQFT}(\overline{\text{Loc}}_1)^{\text{Cauchy}}$$

$$\text{Loc}_1[\text{Cauchy}^{-1}] \simeq BR \Rightarrow \text{time-evolution}$$

$\{ A \in \text{Alg} \text{ w/ } R\text{-action} \}$

$$\boxed{\alpha: R \rightarrow \text{Aut}(A)}$$

Main Theorem of Today [Benini, Giorgetti, AS]

Let $\text{Cloc}_2 := \left\{ \begin{array}{l} \underline{\text{Ob}}: \text{ oriented and time-oriented Lorentzian } 2\text{-mfss } M \\ \text{ that are globally hyperbolic and connected} \\ \underline{\text{Mor}}: \text{ orientation and time-orientation preserving embeddings} \\ f: M \rightarrow M' \text{ s.t. } f(M) \subseteq M' \text{ is causally convex} \\ \text{and } f^*(g') = \Omega^2 g \text{ for some } \Omega^2 \in C^\infty(M, \mathbb{R}^{>0}) \end{array} \right\}$

and $i: C_2^{\text{Diskl}} \xrightarrow{\subseteq} \text{Cloc}_2$ the full subcategory on the two objects

given by the 2d Minkowski spacetime  and the flat cylinder .

Then the inclusion functor i admits a left adjoint

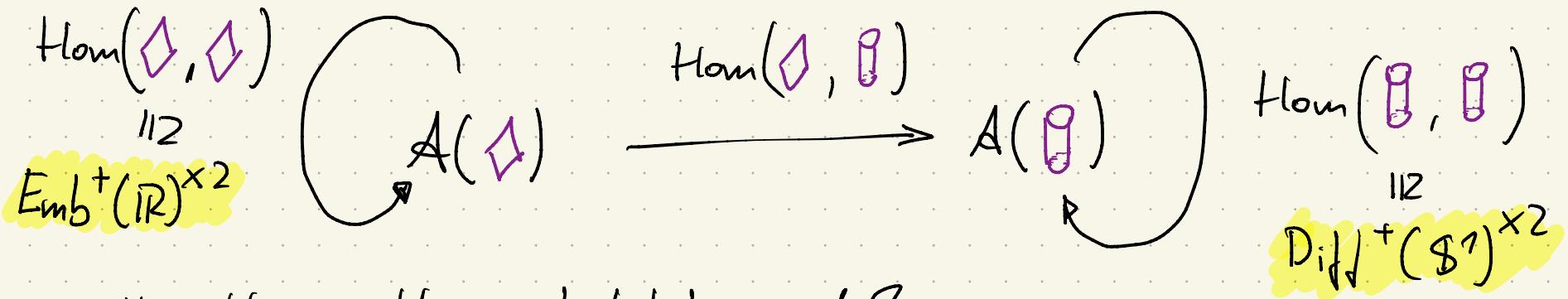
$$L: \text{Cloc}_2 \longrightarrow C_2^{\text{Diskl}}$$

that exhibits C_2^{Diskl} as a localization of Cloc_2 at Causality.

Corollary

Each 2d conformal AQFT satisfying time-slice $A \in \text{AQFT}(\overline{\text{Cloc}_2})$ Cauchy

is completely determined by only two algebras w/ a 1-commutative action:



We call this the skeletal model.

Recon: The skeletal model is easy to compute:

Just restrict the functor $A: \text{Cloc}_2 \rightarrow \text{Alg}$ to $\mathcal{S}^{\text{Disk}} \subseteq \text{Cloc}_2$.

Reconstructing $A \in \text{AQFT}(\overline{\text{Cloc}_2})$ Cauchy from a skeletal model

is much more involved, but can be done via operadic Kan extensions.

Proof sketch for main theorem:

(1) Global conformal embedding theorems [Finster, Müller ; Monclair]

Each $M \in \text{Cloc}_2$ is isomorphic to causally convex open subset $U \subseteq \mathbb{I}$

$\Rightarrow \text{Cloc}_2$ is equivalent to full subcategory $C_2 \subseteq \text{Cloc}_2$ on such U 's

(2) Cloc_2 -morphisms $f: U \subseteq \mathbb{I} \rightarrow U' \subseteq \mathbb{I}$, $(x^+, x^-) \mapsto (f^+(x^+), f^-(x^-))$

extend uniquely to Cauchy developments $Df: DU \subseteq \mathbb{I} \rightarrow DU' \subseteq \mathbb{I}$

\Rightarrow We get a functor $D: C_2 \rightarrow C_2^D := \{U \subseteq \mathbb{I} : DU = U\}$

(3) Because $f: U \rightarrow U'$ is Cauchy iff $Df: DU \rightarrow DU'$ is iso, D is a localization.

(4) $C_2^{\text{Diskl}} \subseteq C_2^D$ is equivalence by a simple conformal geometry argument.

Summing up:

$$\begin{array}{ccccc} & \text{Localization } L \text{ at Cauchy} & & & \\ \text{Cloc}_2 & \xleftarrow[\subseteq]{\sim} & C_2 & \xrightarrow[\subseteq]{D} & C_2^D & \xleftarrow[\subseteq]{\sim} & C_2^{\text{Diskl}} \end{array}$$

□

Application: Chiralization adjunction

- Wanted: Extract chiral components of $A \in \text{AQFT}(\overline{\mathcal{C}\text{loc}_2})$

These should be chiral conformal AQFTS, i.e. theories on

Man_1^{or} w/ $I =$ disjointness.

- Using skeletal models, there exists an obvious candidate:

$$\begin{array}{ccc}
 \pi_{\pm} : & \overline{\mathcal{C}_2^{D,\text{ske}}} & \longrightarrow \overline{\text{Man}_1^{\text{ske}}} \\
 & \begin{array}{c} \diamond, \circlearrowleft \\ \lrcorner, \lrcorner \end{array} & \longrightarrow \begin{array}{c} |, \circlearrowright \\ |, \circlearrowleft \end{array} \\
 ((J^+, J^-) : \lrcorner \rightarrow \lrcorner) & \longrightarrow & (J^\pm : | \rightarrow |) \\
 ([J^+, J^-] : \lrcorner \rightarrow \circlearrowleft) & \longrightarrow & ([J^\pm] : | \rightarrow \circlearrowleft) \\
 ((g^+, g^-) : \lrcorner \rightarrow |) & \longrightarrow & (g^\pm : \circlearrowleft \rightarrow \circlearrowleft)
 \end{array}$$

Thm: [Benini, Giorgatti, AS] Suppose that Alg is complete.

The functor $\Pi_{\pm}^*: \text{AQFT}(\overline{\text{Man}_1^{\text{shl}}}) \longrightarrow \text{AQFT}(\overline{\mathcal{C}_2^{\text{D}, \text{shl}}})$ is

fully faithful and it admits a right adjoint (chiralization functor)

$$\Pi_{\pm *}: \text{AQFT}(\overline{\mathcal{C}_2^{\text{D}, \text{shl}}}) \longrightarrow \text{AQFT}(\overline{\text{Man}_1^{\text{shl}}}).$$

A model for $\Pi_{\pm *}$ is given by taking invariants of the opposite chirality

$$\Pi_{\pm *}(\mathbb{A})(\mathbb{I}) = A(\square) \overset{\text{inv}_{\mp}}{\sim} \begin{array}{l} \text{invariants under } \text{Emb}^+(\mathbb{H}) \\ \text{acting on } x^{\mp} \end{array}$$

$$\Pi_{\pm *}(\mathbb{A})(\mathbb{O}) = A(\square) \overset{\text{inv}_{\mp}}{\sim} \begin{array}{l} \text{invariants under } \text{Diff}^+(\mathbb{S}^1) \\ \text{acting on } x^{\mp} \end{array}$$

Rmk: That's a (locally covariant) generalization and formalization of Rehren's chiral observables.

Example: The Abelian current

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- Solutions on MClo_2 :

$$\text{Sol}(M) = \left\{ j \in \Omega^1(M) : dj = d^*j = 0 \right\}$$

- Linear observables:

$$\mathcal{L}(M) = \frac{\Omega_c^1(M)}{d C_c^\infty(M) \oplus *d C_c^\infty(M)} \cong \frac{\Omega_c^+(M)}{d^+ C_c^\infty(M)} \oplus \frac{\Omega_c^-(M)}{d^- C_c^\infty(M)}$$

$\boxed{\text{self-dual}}$
 $\boxed{\text{anti-self-dual}}$
 $\alpha \alpha = \alpha$ $*\beta = -\beta$

w/ suitable Poisson structure.
 [similar to "F-model" of Fewster, Langy]

- Chiralization of $A = \text{CCR}(\mathfrak{k})$:

$$\pi_\pm(A)(1) \cong \text{CCR}\left(C_c^\infty(\mathbb{R}), \tau_{\text{usual}}\right) \xrightarrow{\sim} \text{usual chiral current}$$

$$\pi_{\pm x}(A)(0) \cong \text{CCR}\left(C^\infty(S^1), \tau_{\text{usual}}\right) \otimes \text{Sym}\left(H_c^\pm(\mathbb{D})\right) \xrightarrow{\sim} \text{topological observables in opposite chirality}$$

(anti-)self-dual cohomology

Topological invariants

- Taking further invariants, we get a functor

$$\text{AQFT}(\overline{\mathcal{C}_2^{\text{D,shL}}}) \xrightarrow{\pi_{\pm}^*} \text{AQFT}(\overline{\text{Man}_1^{\text{shL}}}) \xrightarrow{\text{invariants}} \text{AQFT}(\overline{\mathcal{C}_0})$$

where $\mathcal{C}_0 = (\bullet \rightarrow *)$ w/ $\perp = \{\text{id}_{\bullet}, \text{id}_{\bullet}\}$.

- To $\mathcal{A} \in \text{AQFT}(\overline{\mathcal{C}_2^{\text{D,shL}}})$, this functor assigns the algebra map

$$\begin{array}{ccc} \mathcal{A}(\square) & \xrightarrow{\text{inv}} & \mathcal{A}(\square) \\ \text{commutative} & & \text{not necessarily commutative} \end{array} \quad \begin{array}{l} \text{complete invariants} \\ \text{under } \text{Emb}^+(\mathbb{R})^{+2} \text{ or } \text{Diff}^+(\mathbb{S}^1)^{+2} \end{array}$$

Ex: For the Abelian current:

$$\mathbb{C} \xrightarrow{\text{unit}} \text{Sym}(H_c^1(\square))$$

Question: Do there exist interesting examples with $\mathcal{A}(\square)$ a non-commutative algebra?